

Meson spectrum and tetraquarks through an AdS/QCD inspired potential

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AdS/QCD correspondence allows to compute the quark-antiquark potential in the static limit. We use this piece of information together with the Salpeter equation (Schroedinger equation with relativistic kinematics) and a short range hyperfine splitting potential to describe the bound states of quark-antiquark. Comparison with experiment permits to fit the parameters of the model, namely the quark masses and the quark potential parameters. A discussion of heavy tetraquark masses, prompted by the hypothesis of a diquark-antidiquark structure for the puzzling state X(3872), is also presented.

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A new promising tool to investigate the strong coupling regime of QCD is the AdS/QCD correspondence. In this framework many results have been obtained, namely linear Regge trajectories, the glueball spectrum, the light meson spectrum and decay constants [1]. Through the AdS/QCD correspondence it is also possible to determine the $Q\bar{Q}$ potential [2], the static energy of the strongly interacting quark-antiquark pair.

This result can be used in a quark model to get the meson spectrum [3]; the application can be extended to compute some tetraquark spectra, under the hypothesis that a tetraquark can be considered as a bound state of a diquark and an antidiquark [4].

To compute the meson masses, we solve the semirelativistic wave equation (Salpeter equation):

$$\left(\sqrt{m_1^2 - \nabla^2} + \sqrt{m_2^2 - \nabla^2} + V(r) \right) \psi(\mathbf{r}) = M \psi(\mathbf{r}) \quad (1)$$

where M and ψ are the mass and the wavefunction of the meson, respectively, and m_1 and m_2 are the masses of the constituent quarks. Only the case $\ell = 0$ has been considered. The potential $V(r)$ in (1) comprises three terms:

$$V(r) = V_{AdS/QCD}(r) + V_{spin}(r) + V_0; \quad (2)$$

the last term V_0 is a constant. $V_{AdS/QCD}$ is the AdS/QCD inspired potential, obtained in [2] as a parametric equation:

$$\begin{cases} V_{AdS/QCD}(\lambda) = \frac{g}{\pi} \sqrt{\frac{c}{\lambda}} \left(-1 + \int_0^1 dv v^{-2} \left[e^{\lambda v^2/2} \left(1 - v^4 e^{\lambda(1-v^2)} \right)^{-1/2} - 1 \right] \right) \\ r(\lambda) = 2 \sqrt{\frac{\lambda}{c}} \int_0^1 dv v^2 e^{\lambda(1-v^2)/2} \left(1 - v^4 e^{\lambda(1-v^2)} \right)^{-1/2} \end{cases}, \quad (3)$$

where r is the interquark distance and λ varies in the range: $0 \leq \lambda < 2$. The potential $V_{AdS/QCD}(r)$ depends only on two parameters, g and c . Since the calculation in [2] is performed in the static limit, which corresponds to large quark masses, only heavy meson masses, i.e. of mesons comprising at least one heavy quark, are computed. The second term in (2) takes account of the spin interaction and, in the one-gluon-exchange approximation and in the case $\ell = 0$, can be written as:

$$V_{spin}(r) = A \frac{\delta(r)}{m_1 m_2} (\mathbf{S}_1 \cdot \mathbf{S}_2) \quad \text{with} \quad \delta(r) = \left(\frac{\sigma}{\sqrt{\pi}} \right)^3 e^{-\sigma^2 r^2}. \quad (4)$$

A is a constant proportional to the QCD running coupling constant; to account for the different value of α_s at the two scales, we have introduced two different parameters, A_c in the case of mesons comprising at least one charm quark and A_b for bottom quark.

We have also introduced a cutoff to cure the unphysical divergence at $r = 0$ of the wavefunction solution of the Salpeter equation: at distances shorter than $r_M = k/M$, where k is a constant and M is the mass of the meson, the potential is taken constant, fixed at the value $V(r_M)$. The constant k has different values in the case of mesons made up of identical quarks (k) and mesons made up of different quarks (k').

To solve the Salpeter equation, we use the Multhopp method, a numerical method which allows to transform an integral equation into a set of linear equations, introducing N parameters θ_k , called

Multhopp's angles. The set of equations is (we refer to [5] for further details):

$$\sum_{m=1}^N B_{km} \psi(\theta_k) = M \psi(\theta_m) \quad (5)$$

where $\psi(\theta_k) = u_0(-\cot \theta_k)$ and $\theta_k = \frac{k\pi}{N+1}$ ($k = 1, \dots, N$).

Finally, we fix the parameters of the potential and the constituent quark masses using, as an input, some heavy meson masses. We obtain the following parameters: $c=0.3 \text{ GeV}^2$, $g=2.75$, $V_0=-0.49 \text{ GeV}$, $k=1.48$, $k'=2.15$, $A_c=7.92$, $A_b=3.09$, $\sigma=1.21 \text{ GeV}$, $m_q=0.302 \text{ GeV}$, $m_s=0.454 \text{ GeV}$, $m_c=1.733 \text{ GeV}$, $m_b=5.139 \text{ GeV}$. Using these values, we find the meson spectrum reported in Table 1.

Table 1: Mass spectra for heavy mesons; $q = u, d$. Units are GeV.

Flavor	Level	$J = 0$			$J = 1$		
		Particle	Th. mass	Exp. mass [6]	Particle	Th. mass	Exp. mass [6]
$c\bar{q}$	1S	D	1.862	1.867	D^*	2.027	2.008
	2S		3.393			2.598	
	3S		2.837			2.987	
$c\bar{s}$	1S	D_s	1.973	1.968	D_s^*	2.111	2.112
	2S		2.524			2.670	
	3S		2.958			3.064	
$c\bar{c}$	1S	η_c	2.990	2.980	J/ψ	3.125	3.097
	2S		3.591			3.655	
	3S		3.994			4.047	
$b\bar{q}$	1S	B	5.198	5.279	B^*	5.288	5.325
	2S		5.757			5.819	
	3S		6.176			6.220	
$b\bar{s}$	1S	B_s	5.301	5.366	B_s^*	5.364	5.412
	2S		5.856			5.896	
	3S		6.266			6.296	
$b\bar{c}$	1S	B_c	6.310	6.286	B_c^*	6.338	6.420
	2S		6.869			6.879	
	3S		7.221			7.228	
$b\bar{b}$	1S	η_b	9.387	9.389	Υ	9.405	9.460
	2S		10.036			10.040	
	3S		10.369			10.371	
	4S		10.619			10.620	

The agreement is remarkable, as, for example, in the case of η_b ; in fact, a recent experimental result is [7]:

$$M_{\eta_b} = 9389_{-2.3}^{+3.1} \pm 2.7 \text{ MeV}.$$

This application can be extended to compute tetraquark masses, in order to examine the possibility that some recently observed puzzling states, for example, X(3872) and Y(3940), can be interpreted as diquark-antidiquark bound states. This possibility is justified by the fact that two quarks can attract one another in the one-gluon-exchange approximation, in the $\bar{3}$ color attractive channel. In the same one-gluon-exchange approximation, the energy of this interaction is half of the energy of the interaction between a quark and an antiquark. Following these hints, we have determined the diquark masses, solving the Salpeter equation (1) and using, for the potential $V(r)$,

Table 2: Diquark masses. $\{\}$ denotes a diquark with spin 1 and $[\]$ a diquark with spin 0. Units are GeV.

State	Mass	State	Mass
$\{qs\}$	0.980	$[qs]$	0.979
$\{ss\}$	1.096		
$\{cq\}$	2.168	$[cq]$	2.120
$\{cs\}$	2.276	$[cs]$	2.235
$\{cc\}$	3.414		
$\{bq\}$	5.526	$[bq]$	5.513
$\{bs\}$	5.630	$[bs]$	5.619
$\{bc\}$	6.741	$[bc]$	6.735
$\{bb\}$	10.018		

$1/2 V(r)$ in (2). The results are shown in Table 2, in which the square and the curly brackets indicate a diquark with spin 0 and 1, respectively.

We have computed tetraquark masses again solving Eq. (1), with now m_1 and m_2 the masses of the diquark and the antidiquark constituting the tetraquark. The potential $V(r)$ is the same as the one used for mesons, since the interaction is again between states in the representations $\bar{3}$ and 3 of the SU(3) color group. However, it is necessary to introduce a slight modification, namely an interpolation of the $Q\bar{Q}$ potential (2) with the wavefunctions of the diquark ($\psi_{12}(r)$) and the antidiquark ($\psi_{34}(r)$), due to the fact that diquarks are not point-like particles:

$$\tilde{V}(R) = \frac{1}{N} \int d\mathbf{r}_1 \int d\mathbf{r}_2 |\psi_{12}(\mathbf{r}_1)|^2 |\psi_{34}(\mathbf{r}_2)|^2 V(|\mathbf{R} + \mathbf{r}_1 - \mathbf{r}_2|) \quad (6)$$

with

$$N = \int d\mathbf{r}_1 \int d\mathbf{r}_2 |\psi_{12}(\mathbf{r}_1)|^2 |\psi_{34}(\mathbf{r}_2)|^2. \quad (7)$$

The modified potential $\tilde{V}(r)$ is shown in Fig. 1 in the case of the tetraquark $[cq]\{\bar{c}\bar{q}\}$ (dashed line), together with the AdS/QCD potential (3) (solid line).

In Table 3 the masses computed for the tetraquarks with hidden charm are presented. Let us focus on the 1^{++} state: according to the hypothesis in [4], it should correspond to the particle X(3872), recently observed, for which the identification with a charmonium state is debated [8, 9]. We find $M = 3.899$ GeV, a value not far from experiments. However, confirmation of this interpretation requires the experimental observation of all the other predicted states, for which we obtain the masses reported in Table 3.

Table 3: Four-quark states with hidden charm. Units are GeV.

J^{PC}	Flavor content	Mass (this paper)	Mass [4]	Exp. state	Exp. mass
0^{++}	$[cq][\bar{c}\bar{q}]$	3.857	3.723		
1^{++}	$([cq]\{\bar{c}\bar{q}\} + [\bar{c}\bar{q}]\{cq\})/\sqrt{2}$	3.899	3.872	X(3872)	3.8712 ± 0.0004 [8]
1^{+-}	$([cq]\{\bar{c}\bar{q}\} - [\bar{c}\bar{q}]\{cq\})/\sqrt{2}$	3.899	3.754		
0^{+-}	$\{cq\}\{\bar{c}\bar{q}\}$	3.729	3.832		
1^{+-}	$\{cq\}\{\bar{c}\bar{q}\}$	3.833	3.882		
2^{++}	$\{cq\}\{\bar{c}\bar{q}\}$	3.988	3.952	Y(3940)	$3.943 \pm 0.011 \pm 0.013$ [10]

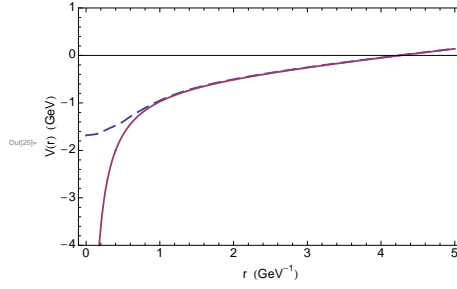


Figure 1: The static energy of a diquark-antidiquark pair (dashed line) and a quark-antiquark pair (solid line). In the first case, the potential corresponds to the $[cq]\{\bar{c}\bar{q}\}$ configuration.

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